

Vessiot's Equivalence Method Applied to Linear Partial Differential Operators

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Outline

- 1 Motivation
- 2 Vessiot Equivalence Method
 - Introductory Example
 - Natural Bundles
 - Symmetry Groupoids
 - Integrability Conditions
 - Equivalence
- 3 Applications to LPDOs
 - Groupoids Θ_q of Gauge Transformations
 - Natural Θ_q -Bundles
 - Invariants
- 4 Examples

Motivation

Example

Linear partial differential operators (LPDOs) of order 2:

$$L = D_{x^1}D_{x^2} + a D_{x^1} + b D_{x^2} + c$$

under gauge transformations:

$$L \mapsto g^{-1}Lg, \quad g = g(x^1, x^2).$$

Questions:

- Equivalence: $L, L' \rightsquigarrow g^{-1}Lg = L'?$
- Invariants?

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Questions:

- Equivalence: $L, L' \rightsquigarrow g^{-1}Lg = L'?$
- Invariants? Laplace:

$$h = c - a_{x^1} - ab, \quad k = c - b_{x^2} - ab.$$

- Generating set of invariants?
- Larger examples?

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Natural bundles

Let X be a manifold, coordinates $(x) = (x^1, \dots, x^n)$.

- $\text{Diff}_{\text{loc}}(X, X)$: local diffeomorphisms $\varphi : X \rightarrow X$.
- A **natural bundle** is a fibre bundle

$$\pi : \mathcal{F} \rightarrow X : (x, u) \rightarrow (x)$$

such that each $\varphi \in \text{Diff}_{\text{loc}}(X, X)$ continues to $\tilde{\varphi} : \mathcal{F} \rightarrow \mathcal{F}$.

- A section $\omega : X \rightarrow \mathcal{F} : (x) \mapsto (x, u = \omega(x))$ is called **geometric object**.

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- φ is an **equivalence** between $\omega, \gamma : X \rightarrow \mathcal{F} \Leftrightarrow$

$$\omega \circ \tilde{\varphi} = \gamma$$

- φ is a **symmetry** of $\omega \Leftrightarrow$

$$\omega \circ \tilde{\varphi} = \omega, \quad \Phi_{\omega(y)}(\varphi, \partial_x \varphi, \dots, \varphi_q) = \omega(x).$$

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- $\psi : \mathcal{F} \rightarrow \mathbb{R}$ is an **invariant** if $\psi \circ \tilde{\varphi} = \psi \quad \forall \varphi \in \text{Diff}_{\text{loc}}(X, X)$.

Symmetry Groupoids

- The jet groupoid $\Pi_q = \Pi_q(X, X)$ has coordinates:

$$(x, y, y_x, \dots, y_q).$$

- Like $\text{Diff}_{\text{loc}}(X, X)$, Π_q acts on \mathcal{F} .
- Symmetry groupoid $\mathcal{R}_q(\omega)$ of $\omega : X \rightarrow \mathcal{F}$ defined by:

$$\Phi_{\omega(y)}(y, y_q) = \omega(x)$$

$$0 \longrightarrow \mathcal{R}_q(\omega) \longrightarrow \Pi_q \begin{array}{c} \xrightarrow{\Phi_\omega} \\ \xrightarrow{\omega} \end{array} \mathcal{F}$$

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- Prolongation

$$0 \longrightarrow \mathcal{R}_{q+1}(\omega) \longrightarrow \Pi_{q+1} \begin{array}{c} \xrightarrow{j_1(\Phi_\omega)} \\ \xrightarrow{j_1(\omega)} \end{array} J_1(\mathcal{F})$$

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Question: Is $\mathcal{R}_q^{(1)}(\omega) = \mathcal{R}_q(\omega)$?

Integrability Conditions

Theorem

$\pi_q^{q+1} : \mathcal{R}_{q+1}(\omega) \rightarrow \mathcal{R}_q(\omega)$ is surjective if and only if there is a section $c : \mathcal{F} \rightarrow \mathcal{F}_1$:

- c is equivariant:

$$c(af_q) = c(a)f_q \quad \forall f_q \in \Pi_q, a \in \mathcal{F}$$

- c fulfills the **Vessiot structure equations**:

$$(I \circ j_1)(\omega) = c(\omega).$$

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If in addition the symbol of $\mathcal{R}_q(\omega)$ is 2-acyclic, $\mathcal{R}_q(\omega)$ is **integrable**, i. e.:

$$\pi_{q+r}^{q+r+s} : \mathcal{R}_{q+r+s}(\omega) \rightarrow \mathcal{R}_{q+r}(\omega) \quad \forall r, s \in \mathbb{N}$$

are all surjective.

Equivalence

Theorem

Two geometric objects ω and γ on \mathcal{F} are equivalent if:

- All invariants ψ on \mathcal{F} coincide for some $x, y \in X$:

$$\psi(\omega(y)) = \psi(\gamma(x)),$$

- $\mathcal{R}_q(\omega), \mathcal{R}_q(\gamma)$ are integrable with $c: \mathcal{F} \rightarrow \mathcal{F}_1$:

$$(I \circ j_1)(\omega) = c(\omega),$$

$$(I \circ j_1)(\gamma) = c(\gamma).$$

Example

In the introductory example (ω, Ω) is equivalent to $(\bar{\omega}, \bar{\Omega})$ if:

$$d\omega = c_1\Omega, \quad d\bar{\omega} = c_1\bar{\Omega}$$

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Groupoids Θ_q of Gauge Transformations

- A simple observation for LPDOs:

$$(g^{-1}Lg) u(x) = g^{-1}L(gu(x))$$

\rightsquigarrow Find a groupoid containing only transformations:

$$(x, u) \mapsto (x, g(x)u).$$

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- On $Y = X \times \mathbb{R}$, coordinates (x, u) , (y, v) define $\Theta_q \leq \Pi_q(Y, Y)$:

$$y = x, \quad u v_u = v.$$

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- All previous slides remain true if Π_q is restricted to Θ_q :

$$0 \longrightarrow \mathcal{R}_q(\omega) \longrightarrow \Theta_q \begin{array}{c} \xrightarrow{\Phi_\omega} \\ \xrightarrow{\omega} \end{array} \mathcal{F}$$

- Natural Θ_q -bundles: Θ_q -action on \mathcal{F} .

Laplace Example – Natural Bundle

- Construct a natural Θ_q -bundle for:

$$L = D_{x^1}D_{x^2} + aD_{x^1} + bD_{x^2} + c.$$

- Gauge transformation:

$$\begin{aligned}g^{-1}Lg &= D_{x^1}D_{x^2} + \left(\frac{g_{x^2}}{g} + a\right)D_{x^1} + \left(\frac{g_{x^1}}{g} + b\right)D_{x^2} \\ &\quad + \frac{g_{x^1x^2}}{g} + a\frac{g_{x^1}}{g} + b\frac{g_{x^2}}{g} + c\end{aligned}$$

- Reminder: $\Theta_q : y = x, \quad u v_u = v.$
 $\Rightarrow v(x, u) = g(x)u.$
- Natural Θ_q -bundle $\mathcal{F} = Y \times \mathbb{R}^3$, coordinates (x, u, a, b, c) :

$$a = \left(\frac{v_{x^2}}{v} + \hat{a}\right),$$

$$b = \left(\frac{v_{x^1}}{v} + \hat{b}\right),$$

$$c = \frac{v_{x^1x^2}}{v} + \hat{a}\frac{v_{x^1}}{v} + \hat{b}\frac{v_{x^2}}{v} + \hat{c}.$$

Laplace Example – Prolongation

$$L = D_{x^1}D_{x^2} + aD_{x^1} + bD_{x^2} + c$$

- Θ_q -action on $\mathcal{F} = Y \times \mathbb{R}^3$:

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- Coordinates of $J_1(\mathcal{F})$:

$$a_{x^1}, \quad a_{x^2}, \quad a_u, \quad b_{x^1}, \quad b_{x^2}, \quad b_u, \quad c_{x^1}, \quad c_{x^2}, \quad c_u.$$

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- Coordinates of $J_{1,X}(\mathcal{F})$:

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Theorem

If the Θ_q -action on \mathcal{F} depends on v_{x^μ}/v only then

- Sections $\omega(x)$ of \mathcal{F} are well-defined.
- $j_1(\omega)(x)$ restricts to the Θ_q -subbundle $J_{1,X}(\mathcal{F}) \subseteq J_1(\mathcal{F})$.

Invariants on Θ_q -bundles

- Θ_q is defined by:

$$y = x, \quad u v_u = v.$$

- Invariant $\psi : J_{r,X}(\mathcal{F}) \rightarrow \mathbb{R}$.
- Derivatives $D_{x^i} \psi$ are invariants on $J_{r+1,X}(\mathcal{F})$.

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- If all coordinates of \mathcal{F}_1 are invariants,

$$\mathcal{F}_1 \cong \mathcal{F} \times \mathbb{R}^k$$

\Rightarrow the invariants on \mathcal{F}_1 are a generating set.

- Next bundle of integrability conditions \mathcal{F}_2 :

$$\mathcal{F}_2 \cong \mathcal{F}_1 \times \mathbb{R}^{nk}$$

Laplace Example – Vessiot Structure Equations

$$L = D_{x^1}D_{x^2} + aD_{x^1} + bD_{x^2} + c$$

- Θ_q -action on $\mathcal{F} = Y \times \mathbb{R}^3$:

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- Coordinates of \mathcal{F}_1 :

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- Better coordinates of \mathcal{F}_1 :

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- Better coordinates of \mathcal{F}_1 :

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- No equivariant sections $\mathcal{F} \rightarrow \mathcal{F}_1$.
- Vessiot structure equations on $\mathcal{F}_2 \cong \mathcal{F}_1 \times \mathbb{R}^4$:

$$h_{x^1} = c_1(h, k), \quad h_{x^2} = c_2(h, k), \quad k_{x^1} = c_3(h, k), \quad k_{x^2} = c_4(h, k).$$

$\Rightarrow \{h, k\}$ is a generating set of invariants.

Example from Shemyakova & Winkler [SW07]

Example

Third order LPDOs under gauge transformations:

$$(D_x + qD_y)D_{xy} + a_{20}D_{xx} + a_{11}D_{xy} + a_{02}D_{yy} + a_{10}D_x + a_{01}D_y + a_{00}$$

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Equivalence conditions: \rightsquigarrow second order invariants.

Order	Σ	Invariants
0	2	$q, I_1 = 2a_{20}q^2 - a_{11}q + 2a_{02}$
1	8	q_x, q_y, I_x^1, I_y^1 4 new invariants
2	15	14 = 6 + 8 old, 1 new
3	21	21 old

Generating set from [SW07] is very compact!

The end.

Thanks!

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